# Exponential functionals of stochastic processes with applications 

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## Outline

(1) Introduction
(2) Disordered systems and Brownian potentials
(3) Exponential functionals in finance

4 Perpetual exponential functionals

Pre-registered goals of the talk:

- Promote interest in exponential functionals and Marc Yor's research.
- Practice some of Beamer's features to build presentation skills.
- Calibrate predicted length of a presentation (plan: 30 minutes).


## Problem 1: disordered systems

In statistical physics, a system with state space $S$ described by Gibbs measures

$$
\mathcal{G}(d x)=\frac{e^{-\beta H(x)} \mathrm{d} x}{\int_{S} e^{-\beta H(y)} \mathrm{d} y} .
$$

Thermodynamic quantities of interest to the physics community tend to rely on computation of the partition function

$$
Z=\int_{S} e^{-\beta H(y)} \mathrm{d} y
$$

or the free energy $\log Z$.
Disordered systems include some inhomogeneity in the Hamiltonian $H$. An example would be random walk vs. the Sinai-Kesten random walk.

## A continuous model

A different related toy model is that of a particle restricted to the interval $[0, L]$ and subject to a random force $F$. To model random inhomogeneities we assume that $F(x)$ is distributed as white noise around some mean value $f_{0}$. Then the potential $U$ is

$$
U(x)=-\int_{0}^{x} F(x) d x=f_{0} x+\sigma B_{x}
$$

with $B$ being a standard Brownian motion. For the sake of simplicity we use scaling properties to get rid of the inverse temperature $\beta$, and index with the drift,

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Z_{L}=\int_{0}^{L} e^{-\beta U(x)} d x
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with $B$ being a standard Brownian motion. For the sake of simplicity we use scaling properties to get rid of the inverse temperature $\beta$, and index with the drift,

$$
Z_{L}^{(\mu, \sigma)}=\int_{0}^{L} e^{-\mu x-\sigma B_{x}} d x
$$

Since thermodynamic variables have the form of expectation values, the partition function must be computed:

$$
g(x)=\frac{\int_{0}^{L} g(x) e^{-\beta U(x)} \mathrm{d} x}{\int_{0}^{L} e^{-\beta U(x)} \mathrm{d} x}
$$

## Special case: zero drift and $L$ exponential random variable

We define a related functional

$$
A_{t}^{(\nu)}=\int_{0}^{t} e^{2\left(\nu s+B_{s}\right)} d s
$$

which can be related to $Z$ through scaling. In particular, we write $A_{t}=A_{t}^{(0)}$. An application of Bougerol's identity

$$
\sinh \left(B_{t}\right) \stackrel{d}{=} W_{A_{t}},
$$

where $W$ Brownian motion independent of the $B$ contained in $A$, allows recovery of much information about the distribution of $A_{t}$. We focus first on the case with $L=L_{\lambda}$ an exponentially distributed random variable with parameter $\lambda>0$.

## Bougerol's identity

Theorem (Alili and Dufresne)
Let $B_{t}, W_{t}$ be two independent standard Brownian motions, and let $A_{t}=\int_{0}^{t} e^{2 B_{s}} \mathrm{~d}$. Then, it holds that

$$
\sinh \left(B_{t}\right) \stackrel{d}{=} W_{A_{t}} .
$$

Sketch of argument.
Consider the Markov process

$$
X_{t}=e^{B_{t}} \int_{0}^{t} e^{-B_{s}} \mathrm{~d} W_{s}, \quad W \perp B
$$

An application of Itô's formula yields the SDE

$$
\mathrm{d} X_{t}=\frac{1}{2} X_{t} \mathrm{~d} t+\left(X_{t} \mathrm{~d} B_{t}+\mathrm{d} W_{t}\right)
$$

The above can equivalently be expressed as

$$
\mathrm{d} X_{t}=\frac{1}{2} X_{t} \mathrm{~d} t+\left(X_{t}^{2}+1\right)^{1 / 2} \mathrm{~d} \beta_{t}
$$

with $\beta_{t}$ another Brownian motion.

## Bougerol's identity

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Then, it holds that

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\sinh \left(B_{t}\right) \stackrel{d}{=} W_{A_{t}} .
$$

Sketch of argument.
An application of Itô's formula to the process $Y_{t}=\sinh \left(\beta_{t}\right)$ yields

$$
\mathrm{d} Y_{t}=\frac{1}{2} Y_{t} \mathrm{~d} t+\left(Y_{t} \mathrm{~d} B_{t}+\mathrm{d} W_{t}\right)
$$

A simple inspection lets us conclude that $X_{t} \stackrel{d}{=} Y_{t}$, for any fixed $t \geq 0$. Finally, conditioning on the process $B$ and applying the Itô isometry allows us to conclude

$$
X_{t} \stackrel{d}{=} \int_{0}^{t} e^{2 B_{s}} \mathrm{~d} s
$$

and the identity now follows.

Special case: zero drift and $L$ exponential random variable
We observe first that by Brownian scaling

$$
\begin{equation*}
\mathbb{E}\left[W_{A_{L_{\lambda}}}^{2 m}\right]=\mathbb{E}\left[N^{2 m}\right] \mathbb{E}\left[A_{L_{\lambda}}^{m}\right], \tag{1}
\end{equation*}
$$

where $N$ is a standard normal r.v.

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where $N$ is a standard normal r.v.
Furthermore, $\left|B_{L_{\lambda}}\right|$ is distributed as an exponential variable with parameter $\theta=\sqrt{2 \lambda}$. As such,

$$
\mathbb{E}\left[\sinh \left(B_{L_{\lambda}}\right)^{2 m}\right]=\int_{0}^{\infty} \theta e^{-\theta x}(\sinh x)^{2 m} \mathrm{~d} x
$$

Provided that $\theta>2 m$, an application of the formula $\sinh x=\frac{1-e^{-2 x}}{2 e^{-x}}$ and a change of variables $x=-\frac{1}{2} \log t$ yields the formula

$$
\begin{equation*}
\mathbb{E}\left[\sinh \left(B_{L_{\lambda}}\right)^{2 m}\right]=\frac{\theta}{2^{2 m}} B\left(\frac{\theta-2 m}{2}, 2 m+1\right) \tag{2}
\end{equation*}
$$

where $B(x, y)=\int_{0}^{1} t^{x-1}(1-t)^{y-1} \mathrm{~d} t$ is the Beta function.

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where $B(x, y)=\int_{0}^{1} t^{x-1}(1-t)^{y-1} \mathrm{~d} t$ is the Beta function.
The use of (1)-(2) with an application of Bougerol's identity $\sinh \left(B_{t}\right) \stackrel{d}{=} W_{A_{t}}$, allows us to obtain the formula, valid for $m \geq 0, \theta>2 m$,

$$
\mathbb{E}\left[A_{L_{\lambda}}^{m}\right]=\frac{\Gamma(1+m) \Gamma(\theta / 2+1) \Gamma(\theta / 2-m)}{2^{m} \Gamma(\theta / 2) \Gamma(1+m+\theta / 2)}
$$

It can be inferred from the formula

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that there are components of a beta and gamma variable, independently of each other.
The following theorem holds.

## Theorem

Let $\lambda>0, a=\sqrt{\lambda / 2}$, and let $L_{\lambda}$ be an exponential variable with parameter $\lambda$, independent of the $B M$. Then we have the identity in distribution

$$
\begin{equation*}
A_{L_{\lambda}} \stackrel{d}{=} \frac{\beta_{1, a}}{2 \Gamma_{a}}, \quad \beta_{1, a} \perp \Gamma_{a}, \tag{3}
\end{equation*}
$$

where $\beta_{1, a}$ denotes a beta variable and $\Gamma_{a}$ a gamma variable, i.e.

$$
\begin{aligned}
\mathbb{P}\left(\beta_{1, a} \in d x\right) & =a(1-x)^{a-1} 1_{(0,1)}(x) d x \\
\mathbb{P}\left(\Gamma_{a} \in d x\right) & =\frac{x^{a-1} e^{-x}}{\Gamma(a)} 1_{(0, \infty)}(x) d x .
\end{aligned}
$$

The identity $A_{L_{\lambda}} \stackrel{d}{=} \frac{\beta_{1, a}}{2 \Gamma_{a}}$ allows us to compute the mean free energy that we were originally interested in. It is the case that

$$
\mathbb{E}\left[\log \Gamma_{a}\right]=\psi(a), \quad \mathbb{E}\left[\log \beta_{1, a}\right]=\psi(1)-\psi(1+a)
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\mathbb{E}\left[\log A_{L_{\lambda}}\right]=-\gamma-\frac{1}{\sqrt{\lambda / 2}}-2 \psi(\sqrt{\lambda / 2})
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If we use Brownian scaling to move back to the original problem, we get the following.

## Proposition

Let a particle be restricted to the interval $[0, L], L$ being exponential with parameter $\lambda>0$, and let the particle be subject to the Brownian potential $U(x)=\sigma B_{x}, B_{x}$ a standard Brownian motion. Then, the partition function and mean free energy of the system is

$$
Z_{L}^{(0, \sigma)}=\int_{0}^{L} e^{-\sigma B_{x}} \mathrm{~d} x, \quad \mathbb{E}\left[\log Z_{L}^{(0, \sigma)}\right]=-2 \log \sigma-\gamma-\frac{1}{\sigma} \sqrt{2 / \lambda}-2 \psi(\sigma \sqrt{\lambda / 2})
$$

Above $\gamma$ denotes the Euler-Mascheroni constant $\gamma:=\lim _{n \rightarrow \infty} \sum_{k=1}^{n} \frac{1}{k}-\log n$.

## What now?

There are several remarks to be made.

- There are several generalizations of Bougerol's identity that allows one to say things about $A_{t}^{\nu}$ for $\nu \neq 0$.


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(1) There are several generalizations of Bougerol's identity that allows one to say things about $A_{t}^{\nu}$ for $\nu \neq 0$. The computations are truly horrifying, sadly.
(2) Marc Yor has proved with another method that there is a more general identity

$$
A_{T_{\lambda}}^{(\nu)} \stackrel{d}{=} \frac{\beta_{1, a}}{2 \Gamma_{b}}, \quad a=\frac{1}{2}\left(\sqrt{\nu^{2}+2 \lambda}+\nu\right), b=\frac{1}{2}\left(\sqrt{\nu^{2}+2 \lambda}-\nu\right) .
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These naturally lead to a whole family of explicitly solvable models with drifted potentials.

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$$

These naturally lead to a whole family of explicitly solvable models with drifted potentials.
(3) If $\nu>0$ and we let $\lambda \rightarrow 0$ in the identity above, we obtain the result

$$
A_{\infty}^{(-\nu)}=\frac{1}{2 \Gamma_{\nu}}
$$

This was first proved by D. Dufresne in an insurance setting.
(9) Knowledge of $\mathbb{E}\left[A_{T_{\lambda}}^{(\nu)}\right]$ also allows one to recover the Laplace transform of $A_{t}$ at fixed $t>0$, provided the exponential variable $T$ is independent of the underlying Brownian motion,

$$
\lambda \int_{0}^{\infty} e^{-\lambda t} \mathbb{E}\left[A_{T_{\lambda}}^{(\nu)} \mid T_{\lambda}=t\right] \mathrm{d} t=\mathbb{E}\left[A_{T_{\lambda}}^{(\nu)}\right]
$$

## References for part 1

A. Comtet, C. Monthus, and M. Yor. Exponential functionals of Brownian motion and disordered systems. Journal of Applied Probability 35, (1998), 255-271.
M. Yor. On Some Exponential Functionals of Brownian Motion. Adv. in Appl. Prob. 24, (1992), 509-531.

- S. Vakeroudi. Bougerol's identity in law and extensions. Probab. Surveys 9, (2012), 411-437.


## Mathematical finance in 2 minutes

- Mathematical finance concerns the dynamics of the value $S_{t}, 0 \leq t \leq T$ of assets, and of computing the distributional properties of $f\left(S_{T}\right)$, for given functions $f$ known as derivatives.
- A typical method is to assume a stochastic setting with asset prices governed by

$$
d S_{t}=r S_{t} d t+\sigma S_{t} d W_{t}
$$

meaning that $S(t)=S_{0} \exp \left(\left(r-\sigma^{2} / 2\right) t+\sigma W_{t}\right)$.

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- For insurance purposes certain types of derivatives have wide use in the industry, especially the class of options that give the contract holder $\max (S(T)-K, 0)=(S(T)-K)^{+}$at the maturity date $T$.
- Financial arguments allow one to conclude that there is a fair price at time $t<T$ for the option

$$
C_{t, T}(K)=e^{-r(T-t)} \mathbb{E}\left[\left(S_{T}-K\right)^{+} \mid \mathcal{F}_{t}\right]
$$

## Problem 2: Asian options

Asian options are also known as Average-Value (AV) options, and for this class of options the payoff function replaces the value $S(T)$ with an average value

$$
\mathcal{A}(T)=\frac{1}{T-t_{0}} \int_{t_{0}}^{T} S_{u} \mathrm{~d} u
$$

where $t_{0}<T$ is the start time of the average value. The corresponding fair price at time $t \in\left[t_{0}, T\right]$,

$$
C_{t, T}\left(K^{\prime}\right)=\frac{e^{-r(T-t)}}{T-t_{0}} \mathbb{E}\left[\left(A_{T}-K^{\prime}\right)^{+} \mid \mathcal{F}_{t}\right], \quad A_{T}=\int_{t_{0}}^{T} S_{u} \mathrm{~d} u
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where $K^{\prime}=K\left(T-t_{0}\right)$.

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Asian option values are normally computed using approximations and numerical techniques due to:

- availability of such methods
- efficiency of modern computational methods
- efficiency of modern hardware
- ...but also due to gaps in the knowledge of the distribution of $A_{t}^{(\nu)}$.
"...it is impossible to derive an analytic expression for an AV-option." (Kemna and Vorst, 1990)
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Recall,

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Assuming $t_{0}<t<T$, it is possible to arrange

$$
A_{T}=\int_{t_{0}}^{T} S_{u} \mathrm{~d} u=\int_{t_{0}}^{t} S_{u} \mathrm{~d} u+S(t) \int_{0}^{T-t} e^{\left(r-\sigma^{2} / 2\right) u+\sigma B_{u}} \mathrm{~d} u
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where $B_{s}=W_{t+s}-W_{t}$. By further factoring out $S(t)$ we obtain a formula in terms of the remaining randomness at time $t$,

$$
C_{t, T}\left(K^{\prime \prime}\right)=\frac{S(t) e^{-r(T-t)}}{T-t_{0}} \mathbb{E}\left[\left(\int_{0}^{T-t} e^{\left(r-\sigma^{2} / 2\right) s+\sigma B_{s}} \mathrm{~d} s-K^{\prime \prime}\right)^{+}\right]
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where we now have $K^{\prime \prime}=\frac{1}{S_{t}}\left(K\left(T-t_{0}\right)-\int_{t_{0}}^{t} S_{u} \mathrm{~d} u\right)$.
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where we now have $K^{\prime \prime}=\frac{1}{S_{t}}\left(K\left(T-t_{0}\right)-\int_{t_{0}}^{t} S_{u} \mathrm{~d} u\right)$.
We note that if $K^{\prime \prime} \leq 0$ the formula has the simplification

$$
\mathbb{E}\left[\left(\int_{0}^{T-t} e^{\left(r-\sigma^{2} / 2\right) s+\sigma B_{s}} \mathrm{~d} s-K^{\prime \prime}\right)^{+}\right]=\mathbb{E}\left[A_{T-t}^{(\nu)}\right]-K^{\prime \prime}, \quad \nu=r-\sigma^{2} / 2
$$

To summarize: assuming the deterministic quantity $\int_{0}^{t_{0}} S_{u} \mathrm{~d} u \geq K\left(T-t_{0}\right)$, we obtain the formula for the Asian option

$$
C_{t, T}(K)=S_{t}\left(\frac{1-e^{-r(T-t)}}{r\left(T-t_{0}\right)}\right)-e^{-(T-t)}\left(K-\frac{1}{T-t_{0}} \int_{t_{0}}^{t} S_{u} \mathrm{~d} u\right)
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Furthermore, we see that knowledge about the distribution of $A_{t}^{\nu}$ is of high interest in financial applications.

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Furthermore, we see that knowledge about the distribution of $A_{t}^{\nu}$ is of high interest in financial applications.

- Further knowledge of the moments $\mathbb{E}\left[\left(A_{t}^{(\nu)}\right)^{m}\right]$ would be useful.
- Knowledge of the distribution of other quantities such as
- $A_{t}^{(\nu)}+x$,
- $\left(A_{t}^{(\nu)}+x\right)^{+}$,
- $\int_{t_{0}}^{T} \exp \left(-X_{s}\right) \mathrm{d} s$, with $X_{t}$ in the class of Lévy processes,
would be highly useful.
Much is known about the distribution of said quantities due to the research of M. Yor, D. Dufresne, and others.

We now let $X$ be a Lévy process and review some results about the distribution of

$$
Z_{t}=\int_{0}^{t} e^{-X_{s}} \mathrm{~d} s
$$

Theorem (Salminen and Vostrikova)
Assume $X_{t}$ is a Lévy process satisfying $\mathbb{E}\left[e^{-\lambda X_{t}}\right]=e^{-t \phi(\lambda)}<\infty$, for all $\lambda \geq 0, t \geq 0$. If further, $\phi(i) \neq \phi(j)$ for all $0 \leq i<j \leq n$, then for $n \in \mathbb{N}$, we have the formula

$$
\mathbb{E}\left[Z_{t}^{n}\right]=n!\sum_{k=0}^{n-1} \frac{e^{-t \phi(k)}-e^{-t \phi(n)}}{\prod_{i=0, i \neq k}^{n}(\phi(i)-\phi(k))} .
$$

The proof relies on the independence of the increments on $X$ and can be carried out by induction.

Some further results when $t=\infty$
In some cases the functional $Z_{\infty}$ is highly tractable for the same kind of analysis.
Corollary (Salminen and Vostrikova)
Let $X$ be a Lévy process with Laplace exponent $\phi$. Define $N:=\min \{n \in \mathbb{N}: \phi(n) \leq 0\}$. Then,

$$
\mathbb{E}\left[Z_{\infty}^{n}\right]= \begin{cases}\frac{n!}{\Pi_{k=1}^{n} \phi(k)}, & \text { if } n<N, \\ +\infty, & \text { if } n \geq N .\end{cases}
$$

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Let $X$ be a Lévy process with Laplace exponent $\phi$. Define $N:=\min \{n \in \mathbb{N}: \phi(n) \leq 0\}$. Then,

$$
\mathbb{E}\left[Z_{\infty}^{n}\right]= \begin{cases}\frac{n!}{\prod_{k=1}^{n} \phi(k)}, & \text { if } n<N, \\ +\infty, & \text { if } n \geq N .\end{cases}
$$

In cases that are not well-behaved for positive integer moments one can sometimes derive information from the negative moments.

Proposition (Bertoin and Yor)
Assume (1) $\mathbb{E}\left[e^{\lambda X_{t}}\right]=e^{t \varphi(\lambda)}<\infty, \forall t, \lambda \geq 0$, and (2) $m:=\mathbb{E}\left[X_{1}\right]=\varphi^{\prime}(0+)>0$. Then, for all $k \in \mathbb{N}$ and $t \in(0,+\infty]$ we have $\mathbb{E}\left[Z_{t}^{-k}\right]<\infty$. The formula

$$
\mathbb{E}\left[Z_{\infty}^{-k}\right]=m \prod_{i=1}^{k-1} \frac{\varphi(i)}{i}
$$

holds. Moreover, if $X_{t}$ lacks positive jumps, then Carleman's Criterion is satisfied and the distribution of $Z_{\infty}$ is determined by its negative integer moments.

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